

EFFECTIVE APPROXIMATION BASED ON BOUNDARY MEASUREMENTS

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An Inverse Multiscale Problem

Objectives

We aim at developing a **methodology to construct coarse approximations** of highly oscillating PDEs when coefficients are *not* known and only *boundary (possibly aggregated) measurements* of its solutions are available. This methodology is **inspired by homogenization theory** but overcome some of its limitations (e.g. periodicity assumptions) and is more **versatile** (e.g. valid outside the homogenization regime).

Assume we are able to measure observables $\mathcal{O}(A_\varepsilon, g)$ associated to the multiscale Schrödinger equation (1) for a few selected loadings g_1, \dots, g_P .

$$\begin{cases} -\operatorname{div}(A_\varepsilon \nabla u_\varepsilon) = 0 & \text{in } \mathcal{D}, \\ (A_\varepsilon \nabla u_\varepsilon) \cdot n = g & \text{on } \partial\mathcal{D}. \end{cases} \quad (1)$$

How can we define an *effective constant coefficient* \bar{A} such that solutions u_ε to (1) for new RHS are well approximated by the solutions \bar{u} to the coarse diffusion problem

$$\begin{cases} -\operatorname{div}(\bar{A} \nabla \bar{u}) = 0 & \text{in } \mathcal{D}, \\ (\bar{A} \nabla \bar{u}) \cdot n = g & \text{on } \partial\mathcal{D}. \end{cases} \quad (2)$$

In this study, the observable is the compliance defined by $\mathcal{O}(A_\varepsilon, g) = \frac{1}{2} \int_{\partial\mathcal{D}} g u_\varepsilon(g)$. Another choice could be boundary measurements $u_\varepsilon(g)|_{\partial\mathcal{D}}$.

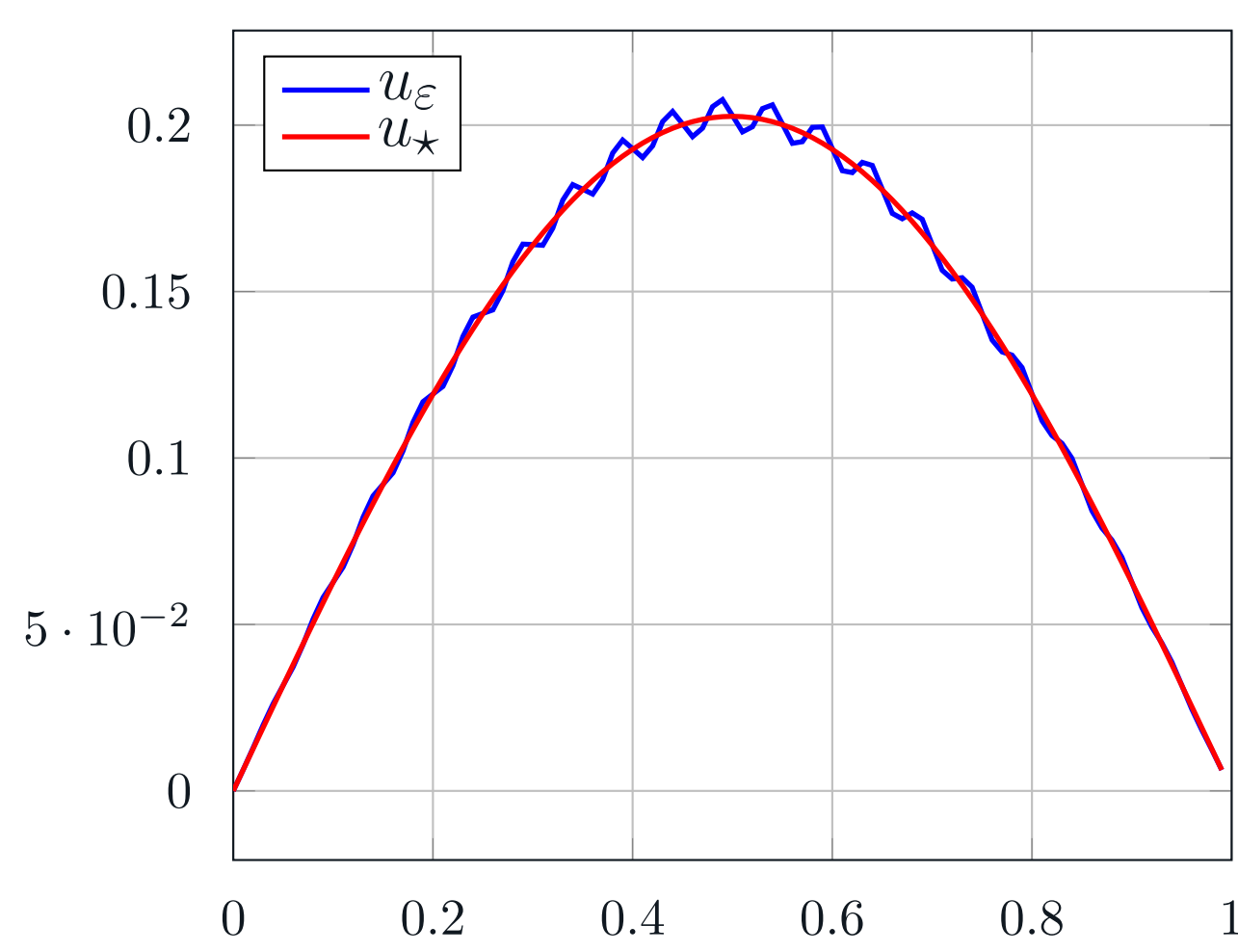


Fig 1. : Diffusion Solutions

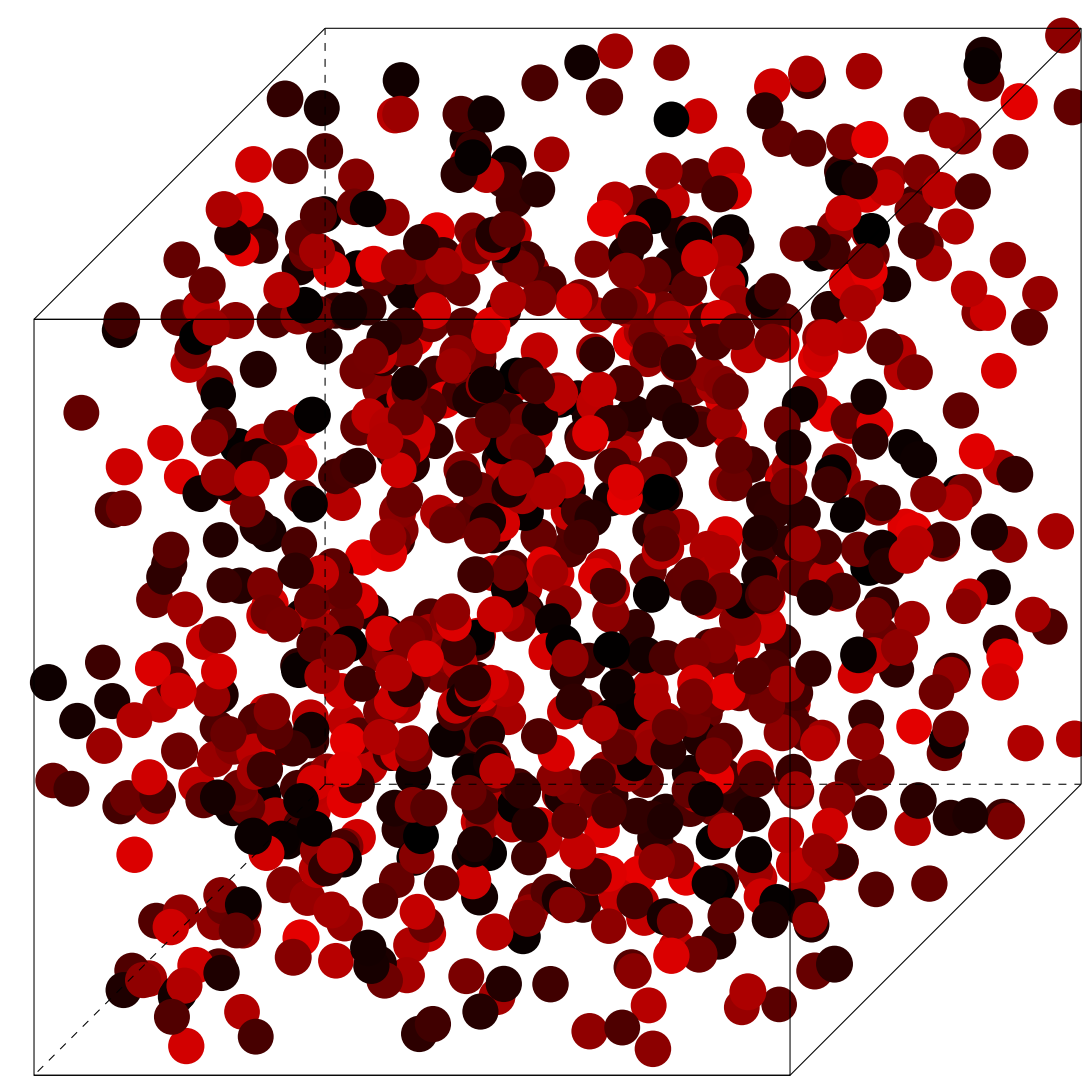


Fig 2. : Multiscale Material

Building An Effective Coefficient

Strategy For Best Effective Coefficient

We examine the **worst case scenario** and try to **minimize** it upon \bar{A} .

$$I_\varepsilon = \inf_{\substack{\bar{A} \in \mathbb{R}_{\text{sym}}^{d \times d} \\ \alpha \leq \bar{A} \leq \beta}} J_\varepsilon(\bar{A}),$$

with

$$J_\varepsilon(\bar{A}) = \sup_{\substack{g \in L^2(\partial\mathcal{D}), \\ \|g\|_{L^2(\partial\mathcal{D})} = 1}} \left(|\mathcal{O}(A_\varepsilon, g) - \mathcal{O}(\bar{A}, g)| \right).$$

The arginf $\bar{A}_\varepsilon^{\text{opt}} \in \mathbb{R}_{\text{sym}}^{d \times d}$ gives a satisfying effective description of the system : **the related solutions \bar{u} are good L^2 approximation of u_ε .**

Numerical Aspects :

- Cost Function** : The sup is replaced by a max over the space spanned by a few loadings g_1, \dots, g_P (in practice $P \approx 3$). -
- Computational Cost** : Each step requires solving a P coarse PDE. Using a gradient descent with adaptative step size (e.g. Armijo Rule) We need $N_{\text{iter}} \approx 15$ iterations.

Comparison To Other Strategies

Our strategy is consistent with Homogenization theory (see [1]) in the sense that

Consistency with Homogenization Theory

In the **periodic setting** (i.e. $A_\varepsilon(x) = A_{\text{per}}(\frac{x}{\varepsilon})$), any sequence of *quasiminimizers* $(\bar{A}_\varepsilon^\#)_{\varepsilon>0}$ that satisfies $I_\varepsilon \leq J_\varepsilon(\bar{A}_\varepsilon^\#) \leq I_\varepsilon + \text{err}(\varepsilon)$ converges to the homogenized matrix A_\star :

$$\lim_{\varepsilon \rightarrow 0} \bar{A}_\varepsilon^\# = A_\star.$$

In [2], another constant effective coefficient $\bar{B}_\varepsilon^{\text{opt}} \in \mathbb{R}_{\text{sym}}^{d \times d}$ is constructed using the knowledge of *full field measurements* $u_\varepsilon(g)$: $\bar{B}_\varepsilon^{\text{opt}} = \operatorname{arginf}_{\bar{A} \in \mathbb{R}_{\text{sym}}^{d \times d}} \sup_{g \in L^2(\partial\mathcal{D})} \|u_\varepsilon(g) - u(\bar{A}, g)\|_{L^2(\mathcal{D})} / \|g\|_{L^2(\partial\mathcal{D})}$.

We can compare the performance of both coefficients $\bar{A}_\varepsilon^{\text{opt}}$ and $\bar{B}_\varepsilon^{\text{opt}}$ using the criteria

$$\operatorname{Err}(\bar{A}) = \sup_{g \in L^2(\partial\mathcal{D})} \frac{\|u_\varepsilon(g) - u(\bar{A}, g)\|_{L^2(\mathcal{D})}}{\|u_\varepsilon(g)\|_{L^2(\mathcal{D})}}.$$

Using Complete Field Measurements

There exists $C_1, C_2 > 0$ such that for any $\varepsilon > 0$,

$$\begin{aligned} \operatorname{Err}(\bar{B}_\varepsilon^{\text{opt}}) &\leq C_1 \operatorname{Err}(\bar{A}_\varepsilon^{\text{opt}}), \\ \operatorname{Err}(\bar{A}_\varepsilon^{\text{opt}}) &\leq C_2 \sqrt{\operatorname{Err}(\bar{B}_\varepsilon^{\text{opt}})}. \end{aligned}$$

References

- [1] A. Bensoussan, J.-L. Lions, and G. Papanicolaou. *Asymptotic analysis for periodic structures*. Vol. 374. American Mathematical Soc., 2011.
- [2] C. Le Bris, F. Legoll, and S. Lemaire. "On the best constant matrix approximating an oscillatory matrix-valued coefficient in divergence-form operators". In: *ESAIM: Control, Optimisation and Calculus of Variations* 24.4 (2018), pp. 1345–1380.

Numerical Results

Experiments have been conducted in 2D ($\mathcal{D} = [0, 1]^2$) using the coefficient

$$A_{\text{per}}(x, y) = \begin{pmatrix} 22 + 10(\sin(2\pi x) + \sin(2\pi y)) & 0 \\ 0 & 12 + 2(\sin(2\pi x) + \sin(2\pi y)) \end{pmatrix},$$

and the loadings $g_{p,q}(x, y) = \cos(p\pi x) \cos(q\pi y)$.

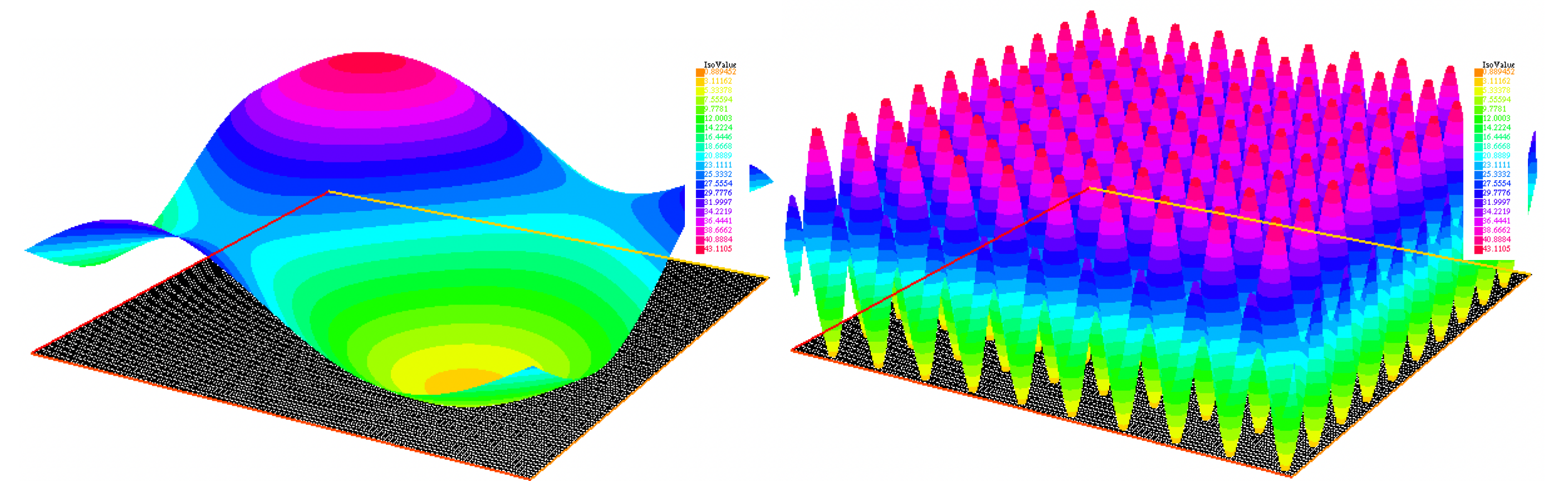


Fig 3 : Kernel Coefficient $A_{\text{per},11}$ and Oscillating Coefficient $A_{\varepsilon,11} = A_{\text{per},11}(\frac{\cdot}{\varepsilon})$ (with $\varepsilon = 0.1$).

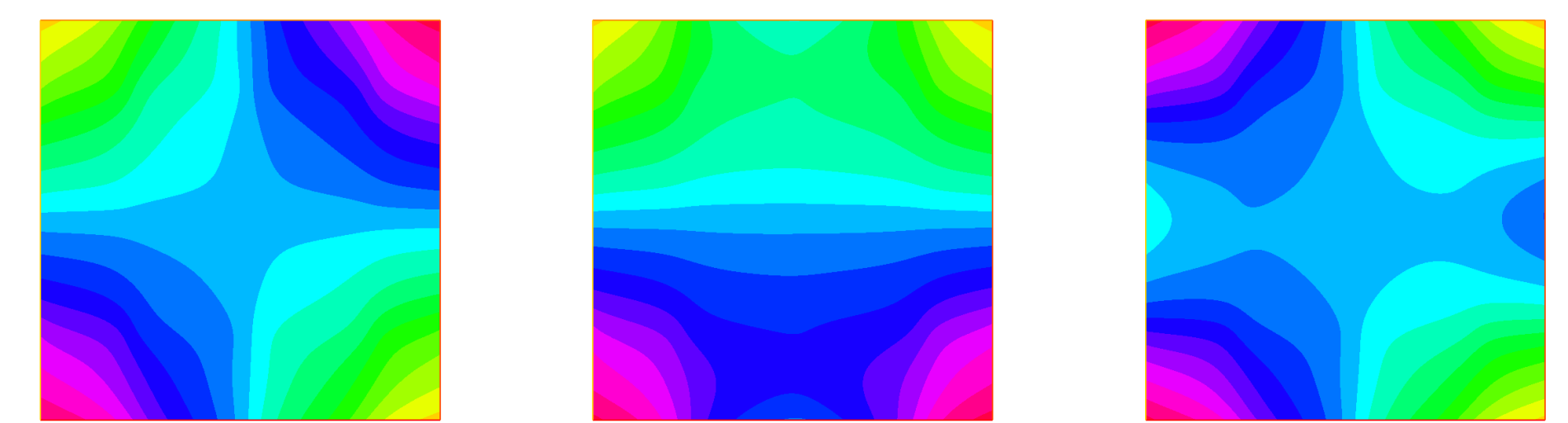


Fig 4 : Loadings $g_{p,q}$.

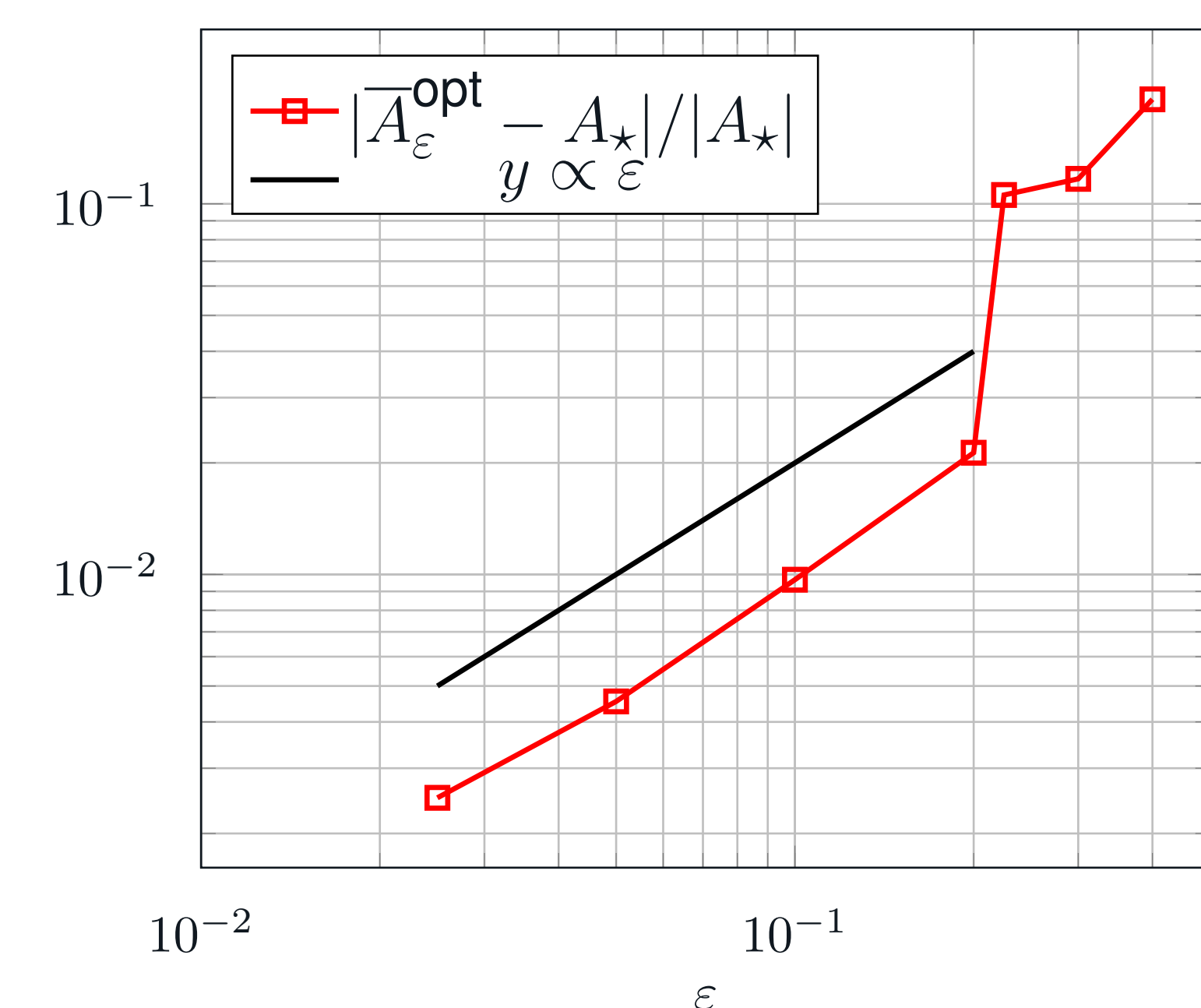


Fig 5 : Error between homogenized coefficient A_\star and effective coefficient $\bar{A}_\varepsilon^{\text{opt}}$ as a function of ε .

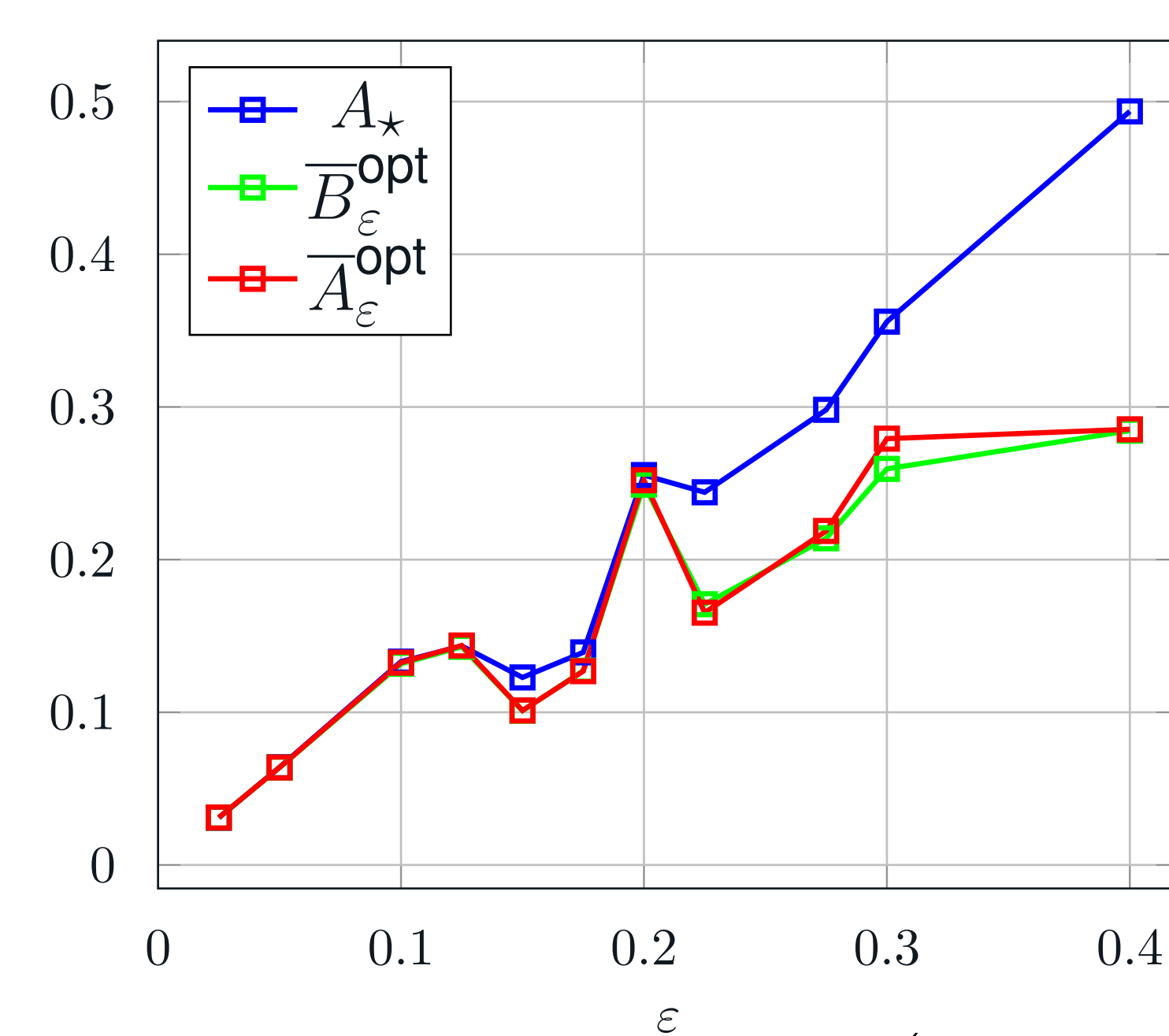


Fig 6 : L^2 maximal error $\operatorname{Err}_\varepsilon(\bar{A}) = \sup_{f \in \operatorname{Span}_{1 \leq p \leq 11}(g_p)} \left(\frac{\|u_\varepsilon(g) - \bar{u}(\bar{A}, g)\|_{L^2(\Omega)}}{\|u_\varepsilon(g)\|_{L^2(\Omega)}} \right)$ for various constant coefficient \bar{A} ($\bar{A}_\varepsilon^{\text{opt}}$ and $\bar{B}_\varepsilon^{\text{opt}}$ computed using only g_1, g_2 and g_3).

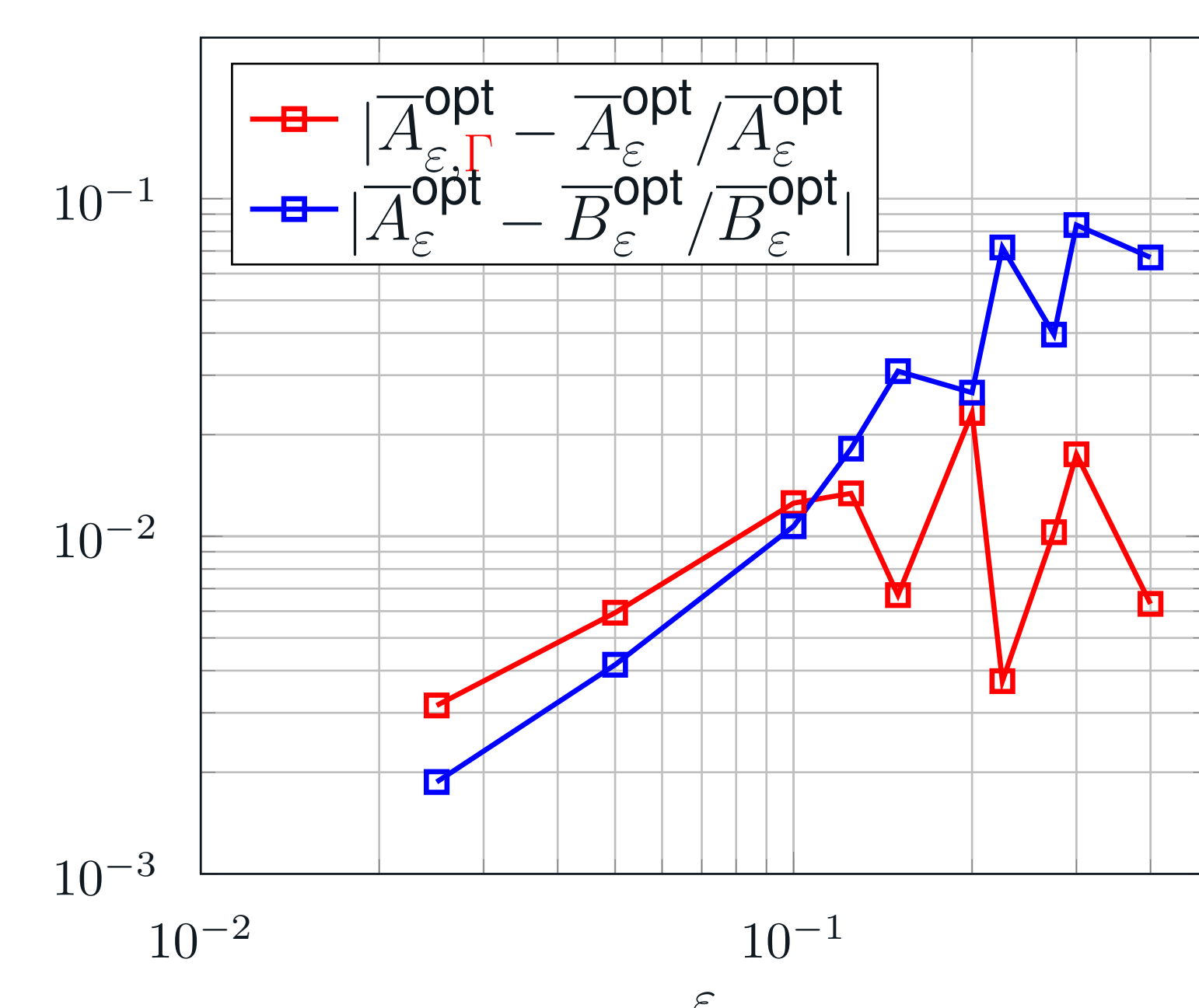


Fig 7 : Relative error between $\bar{A}_\varepsilon^{\text{opt}}$ and $\bar{A}_{\varepsilon,\Gamma}^{\text{opt}}$ (constructed using partial observables $\mathcal{O}_\Gamma(A_\varepsilon, g) = \int_\Gamma g u_\varepsilon(g)$ with $\Gamma \subsetneq \partial\mathcal{D}$ and $|\Gamma| = 0.5|\partial\mathcal{D}|$).

